Implementation of the First Collision Source Method in a Three-Dimensional, Unstructured Tetrahedral Mesh, Discrete-Ordinates Code

Jong Woon Kim*, Ser Gi Hong, and Young-Ouk Lee Korea Atomic Energy Research Institute, 1045 Daedeokdaero, Yuseong-gu, Daejeon, Korea, 305-353 *Corresponding author: jwkim@kaeri.re.kr

1. Introduction

As with any discrete-ordinates code, ray-effects are an inherent problem, especially for shielding type problems with optically thin regions and localized (point) sources. Although ray-effects may be mitigated by increasing the quadrature order, this is often computationally prohibitive.

To mitigate ray-effects, many discrete-ordinates codes use first collision source methods. Such methods are characterized by a decomposition of the flux into its uncollided and collided components.

The uncollided flux is calculated analytically and the collided flux is calculated with the discrete-ordinates method.

In this paper, the implementation of the first collision source method in the MUST (Multi-group Unstructured geometry S_N Transport) [1] code is presented and the results are compared with that of Attila [2,3].

2. Method and Results

2.1 Derivation of the Governing Equation

The one-group transport equation [4] with vacuum boundaries and an isotropic point source at \vec{r}_{p} is

$$\vec{\Omega} \cdot \nabla \psi(\vec{r}, \vec{\Omega}) + \sigma_{t}(\vec{r}) \psi(\vec{r}, \vec{\Omega})
= \sum_{\ell=0}^{L} \sigma_{s,\ell}(\vec{r}) \sum_{m=-\ell}^{\ell} Y_{\ell,m}^{*}(\vec{\Omega}) \Phi_{\ell}^{m}(\vec{r}) + \frac{q_{p}}{4\pi} \delta(\vec{r} - \vec{r}_{p}), \tag{1}$$

where the moments of the angular flux is

$$\Phi_{\ell}^{m}(\vec{r}) = \int d\vec{\Omega} Y_{\ell m}(\vec{\Omega}) \psi(\vec{r}, \vec{\Omega}). \tag{2}$$

Decompose the angular flux into two components as

$$\psi(\vec{r}, \vec{\Omega}) = \psi^{(u)}(\vec{r}, \vec{\Omega}) + \psi^{(c)}(\vec{r}, \vec{\Omega}), \tag{3}$$

where $\psi^{(u)}$ is the uncollided angular flux and $\psi^{(c)}$ is the collided angular flux.

Rewrite Eq. (1) with Eq. (3) then, we have two equations

$$\vec{\Omega} \cdot \nabla \psi^{(u)}(\vec{r}, \vec{\Omega}) + \sigma_t(\vec{r}) \psi^{(u)}(\vec{r}, \vec{\Omega}) = \frac{q_p}{4\pi} \delta(\vec{r} - \vec{r}_p), \tag{4a}$$

$$\vec{\Omega} \bullet \nabla \psi^{(c)}(\vec{r}, \vec{\Omega}) + \sigma_{\bullet}(\vec{r}) \psi^{(c)}(\vec{r}, \vec{\Omega})$$

$$= \sum_{\ell=0}^{L} \sigma_{s,\ell}(\vec{r}) \sum_{m=-\ell}^{\ell} Y_{\ell,m}^{*}(\vec{\Omega}) \Phi_{\ell}^{m,(c)}(\vec{r}) + q_{s}^{(u)}(\vec{r}).$$
 (4b)

Eq. (4a) can be solved analytically for uncollided angular flux and r_{\min} is introduced to prevent it from dividing by zero or very small distance.

$$\psi^{(u)}(\vec{r}, \vec{\Omega}) = \begin{cases} \delta(\vec{\Omega} - \vec{\Omega}_r) \frac{q_p}{4\pi} \frac{e^{-\tau(\vec{r}, \vec{r}_p)}}{\left|\vec{r} - \vec{r}_p\right|^2}, & |\vec{r} - \vec{r}_p| \ge r_{\min}, \\ \delta(\vec{\Omega} - \vec{\Omega}_r) \frac{q_p}{4\pi} \frac{1}{r_{\min}^2}, & |\vec{r} - \vec{r}_p| \le r_{\min}. \end{cases}$$
(5)

The first collision source, $q_s^{(u)}$, is defined as

$$q_{s}^{(u)}(\vec{r}) = \sum_{\ell=0}^{L} \sigma_{s,\ell}(\vec{r}) \sum_{m=-\ell}^{\ell} Y_{\ell,m}^{*}(\vec{\Omega}) \Phi_{\ell}^{m,(u)}(\vec{r}).$$
 (6)

The spherical harmonic moments of the uncollided angular flux, $\Phi_\ell^{m,(u)}$, is calculated as

$$\begin{split} \Phi_{\ell}^{m,(u)} &= \int d\overrightarrow{\Omega} Y_{\ell,m}(\overrightarrow{\Omega}) \psi^{(u)}(\overrightarrow{r},\overrightarrow{\Omega}) \\ &= \begin{cases} Y_{\ell,m}(\overrightarrow{\Omega}_r) \frac{q_p}{4\pi} \frac{e^{-\tau(\overrightarrow{r},\overrightarrow{r}_p)}}{\left|\overrightarrow{r}-\overrightarrow{r}_p\right|^2}, & \left|\overrightarrow{r}-\overrightarrow{r}_p\right| \geq r_{\min}, \\ Y_{\ell,m}(\overrightarrow{\Omega}_r) \frac{q_p}{4\pi} \frac{1}{r_{\min}^2}, & \left|\overrightarrow{r}-\overrightarrow{r}_p\right| < r_{\min}, \end{cases} \end{split} \tag{7}$$

where $\tau(\vec{r}, \vec{r}_p)$ is the optical distance between \vec{r} and \vec{r}_p .

The optical distance from the point source to each node is calculated with

$$\tau(\vec{r}, \vec{r}_p) = \sum_{i=1}^{N} r_i \sigma_{t,i}.$$
point
source
$$mat1$$

$$r_p$$

Fig. 1. The optical distance at the material interface.

2.2 Test and Results

For verification, the test problem is configured as shown in Fig. 2 (point source is located at the corner of 10cm cube).

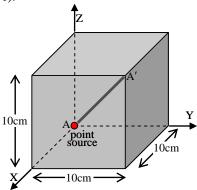


Fig. 2. The configuration of the test problem.

The unstructured tetrahedral mesh is generated by TetGen [5] code. The cross sections and parameters are listed in Table I. The reference scalar flux along the AA' line is calculated by Attila with the identical conditions.

To compare the efficacy of the first collision source method, another run with very small volume source (identical source strength; 10#/sec) is performed with high S_N order (S_{16}) and the result is shown in Fig. 3.

Table I: Cross sections and parameters

	With first collision	Without first collision
# of cell	Attila: 39647 MUST: 39742	MUST: 39837
$\Sigma_{\text{tot}} (\text{cm}^{-1})$	9.3e-2	
$\Sigma_{\rm s} ({\rm cm}^{-1})$	1.0e-2	
Point source	10.0	
(#/sec)	at (0.05, 0.05, 0.05)	-
Volume source		10000.0
(#/cm ³ -sec)	-	$(0.1 \times 0.1 \times 0.1 \text{cm}^3)$
S_N order	4	16
Error criterion	1.0e-11	
r_{min}	0.01	-

In Fig. 3, we can see 36 high-flux points on the opposite side of the source point location. These are the typical ray-effects phenomenon when the discrete ordinate calculation is performed with localized source in spite of using high S_N order.

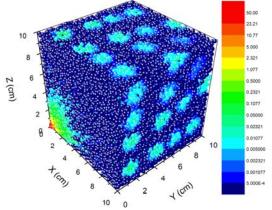


Fig. 3. Scalar flux without a first collision source method (S_{16}) .

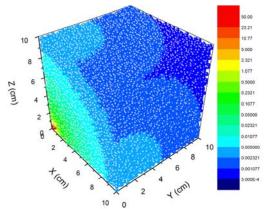


Fig. 4. Scalar flux with a first collision source method (S₄).

However, using the first collision source method, a very low S_N order (even S_4) can resolve the solution, as can be seen from Fig. 4.

To check the calculation result, the scalar flux along AA' line is compared with the result of Attila and shown in Fig. 5.

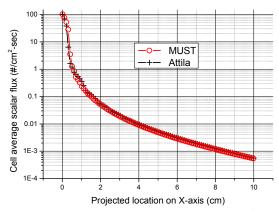


Fig. 5. The comparison of scalar fluxes along AA' line.

The two scalar fluxes match well each other and we confirmed that the first collision source method is successfully implemented in the MUST code.

3. Conclusions

The first collision source method is widely used to mitigate ray-effects in the discrete ordinates code.

The typical ray-effects phenomenon is exampled by using the 3D simple test problem and the scalar flux distributions that are calculated with/without the first collision source method are compared.

The numerical test shows that the first collision source method which is implemented in the MUST code gives accurate solution.

Acknowledgement

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